

A tale of two metals: How strong interactions can destroy an ordinary metal.

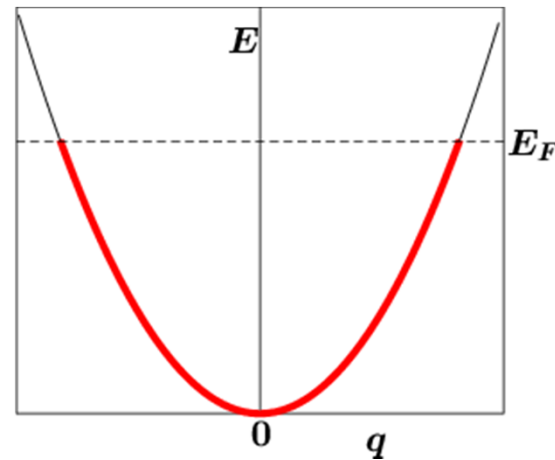
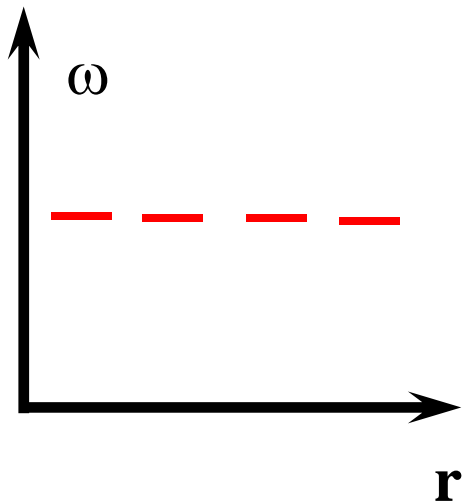
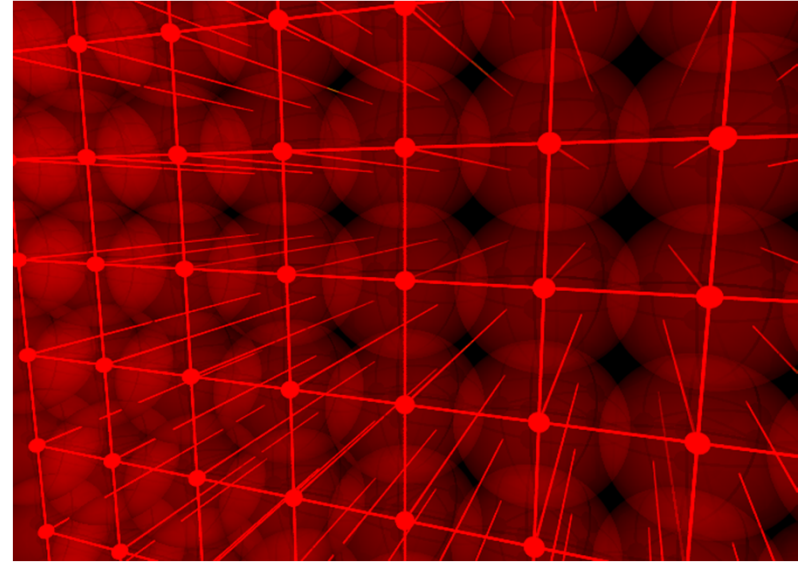
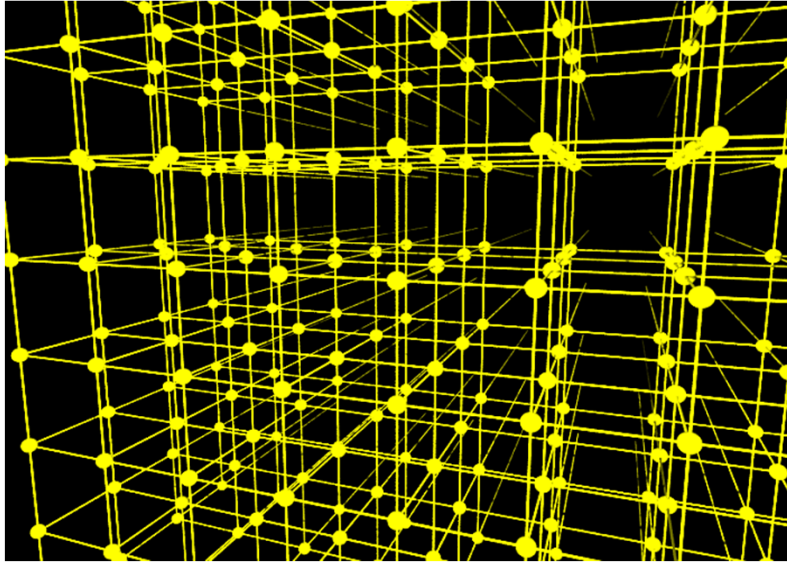
A.-M. Tremblay



Theory Canada 8, Bishop's, 2013



How to make a metal



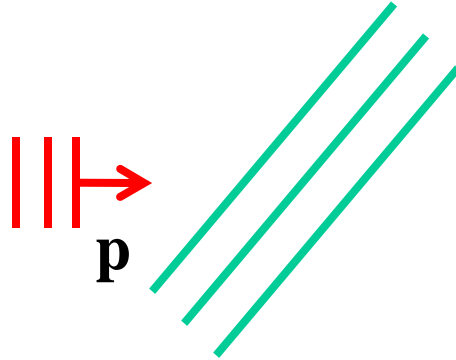
Courtesy, S. Julian



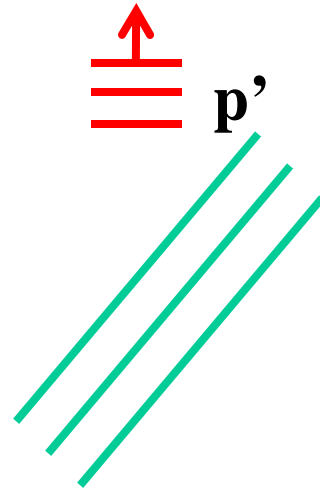
Superconductivity



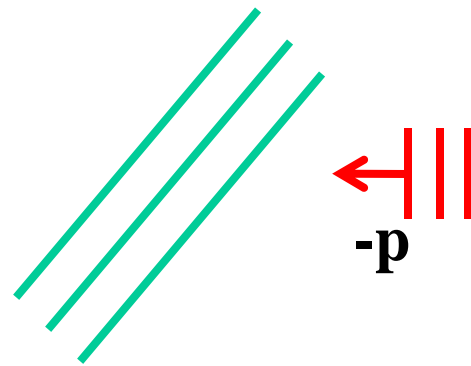
Attraction mechanism in the metallic state



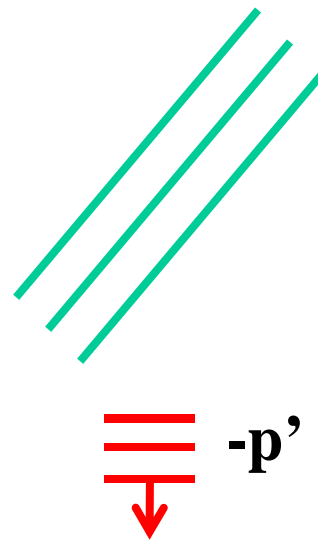
Attraction mechanism in the metallic state



Attraction mechanism in the metallic state



Attraction mechanism in the metallic state



#1 Cooper pair, #2 Phase coherence

$$E_P = \sum_{\mathbf{p}, \mathbf{p}'} U_{\mathbf{p}-\mathbf{p}'} \psi_{\mathbf{p}\uparrow, -\mathbf{p}\downarrow} \psi_{\mathbf{p}'\uparrow, -\mathbf{p}'\downarrow}^*$$

$$E_P = \sum_{\mathbf{p}, \mathbf{p}'} U_{\mathbf{p}-\mathbf{p}'} \left(\langle \psi_{\mathbf{p}\uparrow, -\mathbf{p}\downarrow} \rangle \psi_{\mathbf{p}'\uparrow, -\mathbf{p}'\downarrow}^* + \psi_{\mathbf{p}\uparrow, -\mathbf{p}\downarrow} \langle \psi_{\mathbf{p}'\uparrow, -\mathbf{p}'\downarrow}^* \rangle \right)$$

$$|\text{BCS}(\theta)\rangle = \dots + e^{iN\theta} |N\rangle + e^{i(N+2)\theta} |N+2\rangle + \dots$$



Half-filled band is metallic?

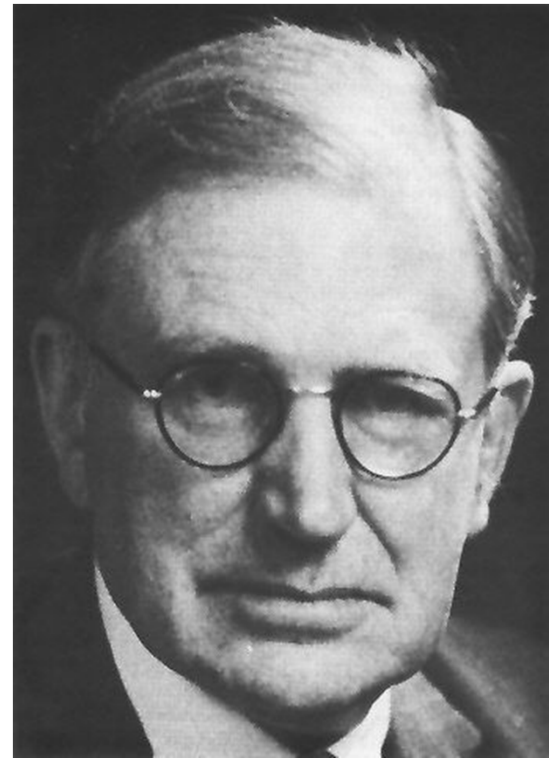


Half-filled band: Not always a metal

NiO, Boer and Verwey



Peierls, 1937



Mott, 1949



« Conventional » Mott transition

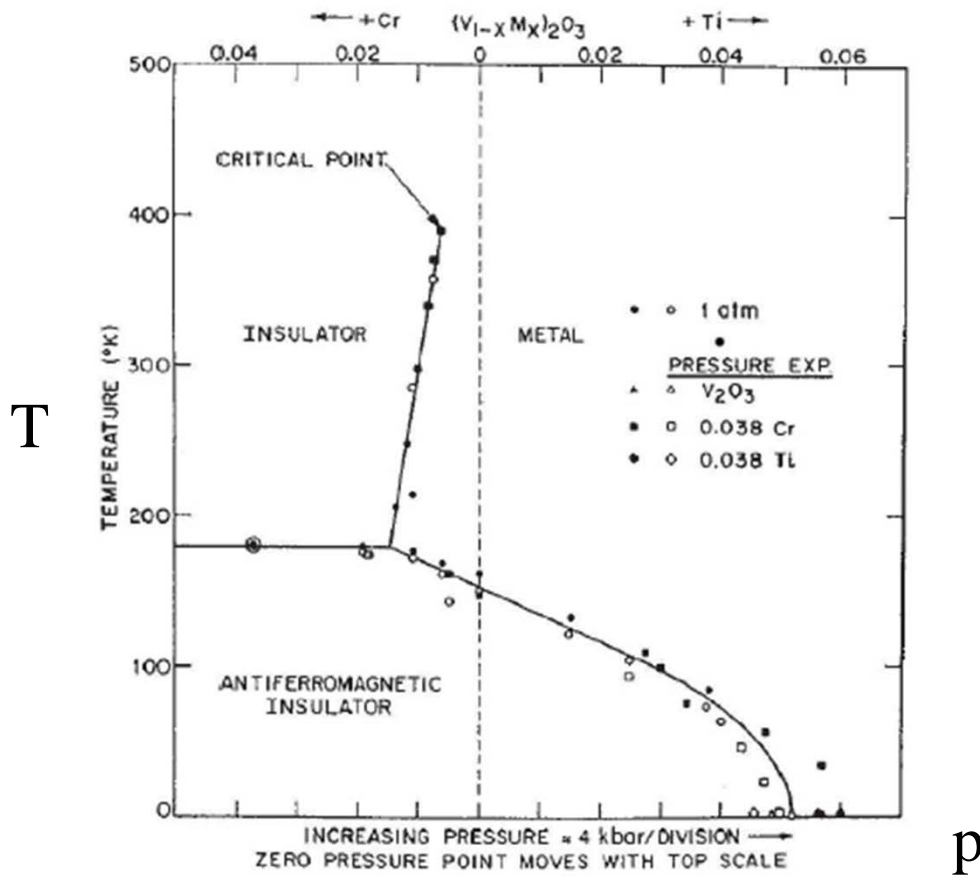
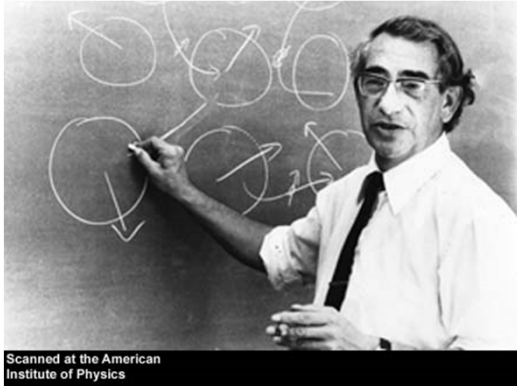
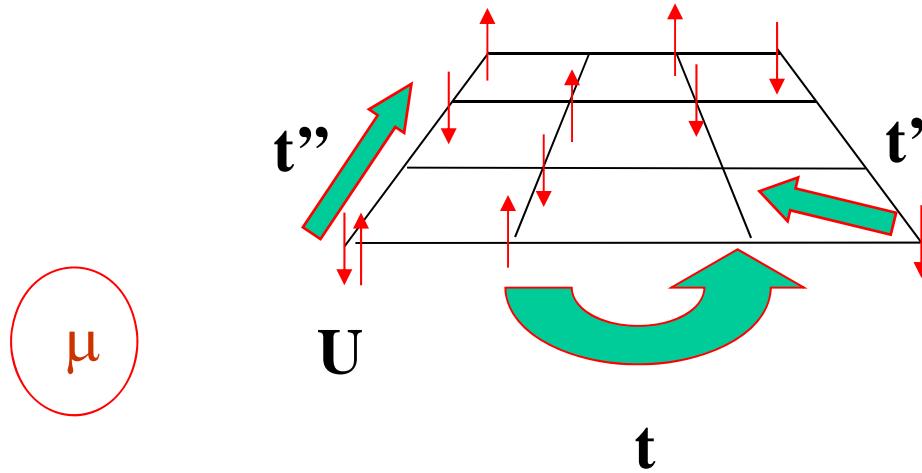


Figure: McWhan, PRB 1970; Limelette, Science 2003

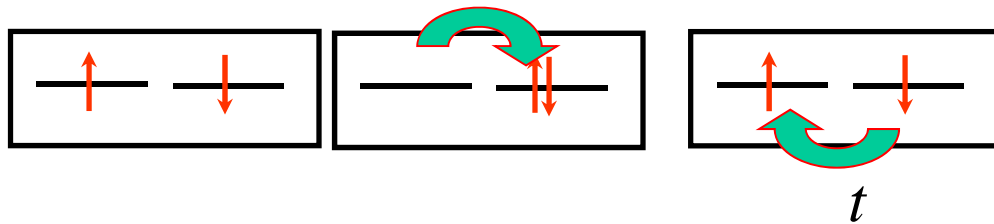
Hubbard model



1931-1980



$$H = - \sum_{\langle ij \rangle \sigma} t_{i,j} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

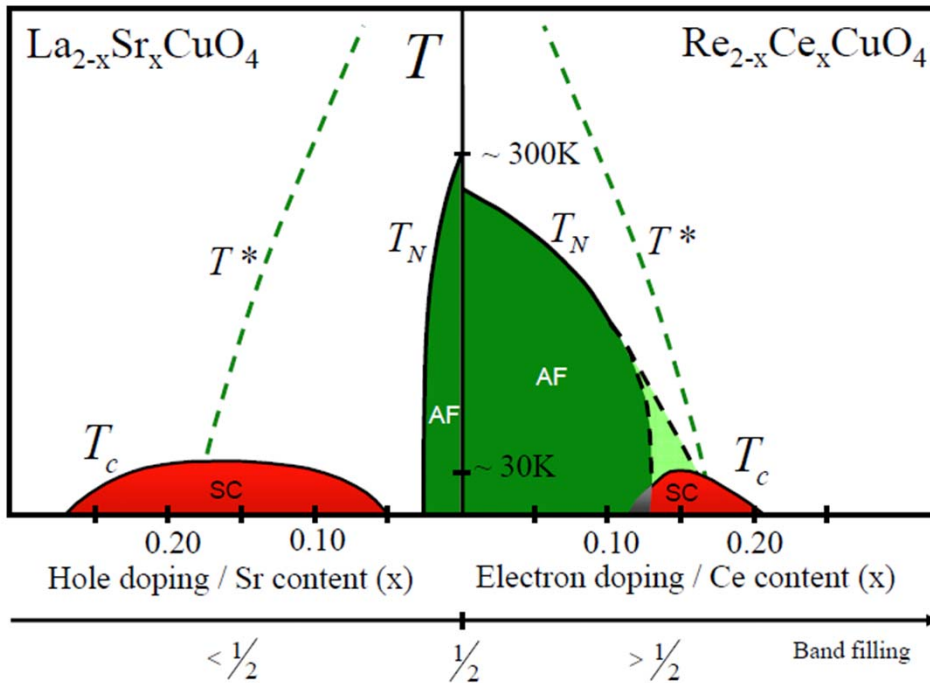


Effective model, Heisenberg: $J = 4t^2 / U$



High-temperature superconductors

Armitage, Fournier, Greene, RMP (2009)



- Competing order

- Current loops: Varma, PRB **81**, 064515 (2010)
- Stripes or nematic: Kivelson et al. RMP **75** 1201(2003); J.C.Davis
- d-density wave : Chakravarty, Nayak, Phys. Rev. B **63**, 094503 (2001); Affleck et al. flux phase
- SDW: Sachdev PRB **80**, 155129 (2009) ...

- Or Mott Physics?

- RVB: P.A. Lee Rep. Prog. Phys. **71**, 012501 (2008)

What is under the dome?
Mott Physics away from $n = 1$

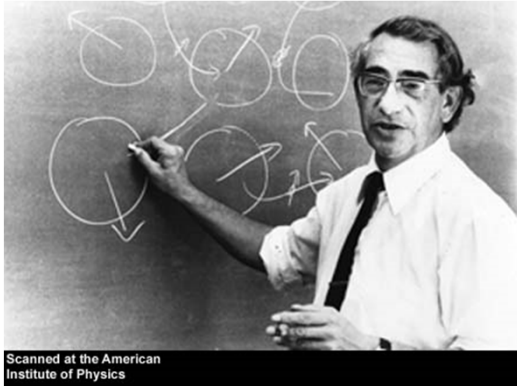


Outline

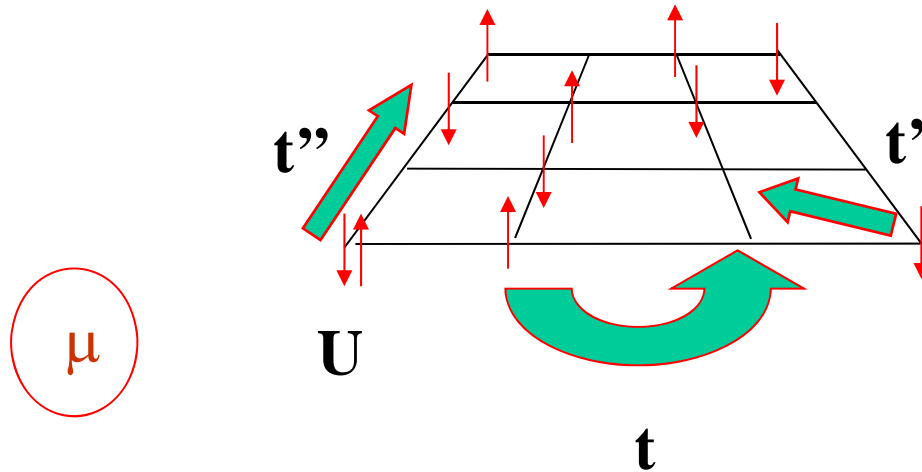
- Pseudogap
 - From AFM fluctuations
- **Weak vs strong coupling**
- Dynamical mean-field theory: crossover
- Phase diagram and pseudogap at strong coupling



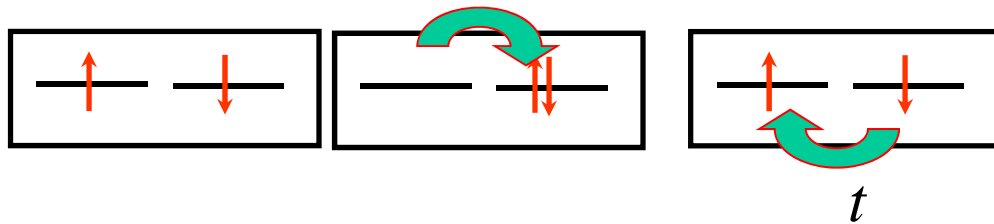
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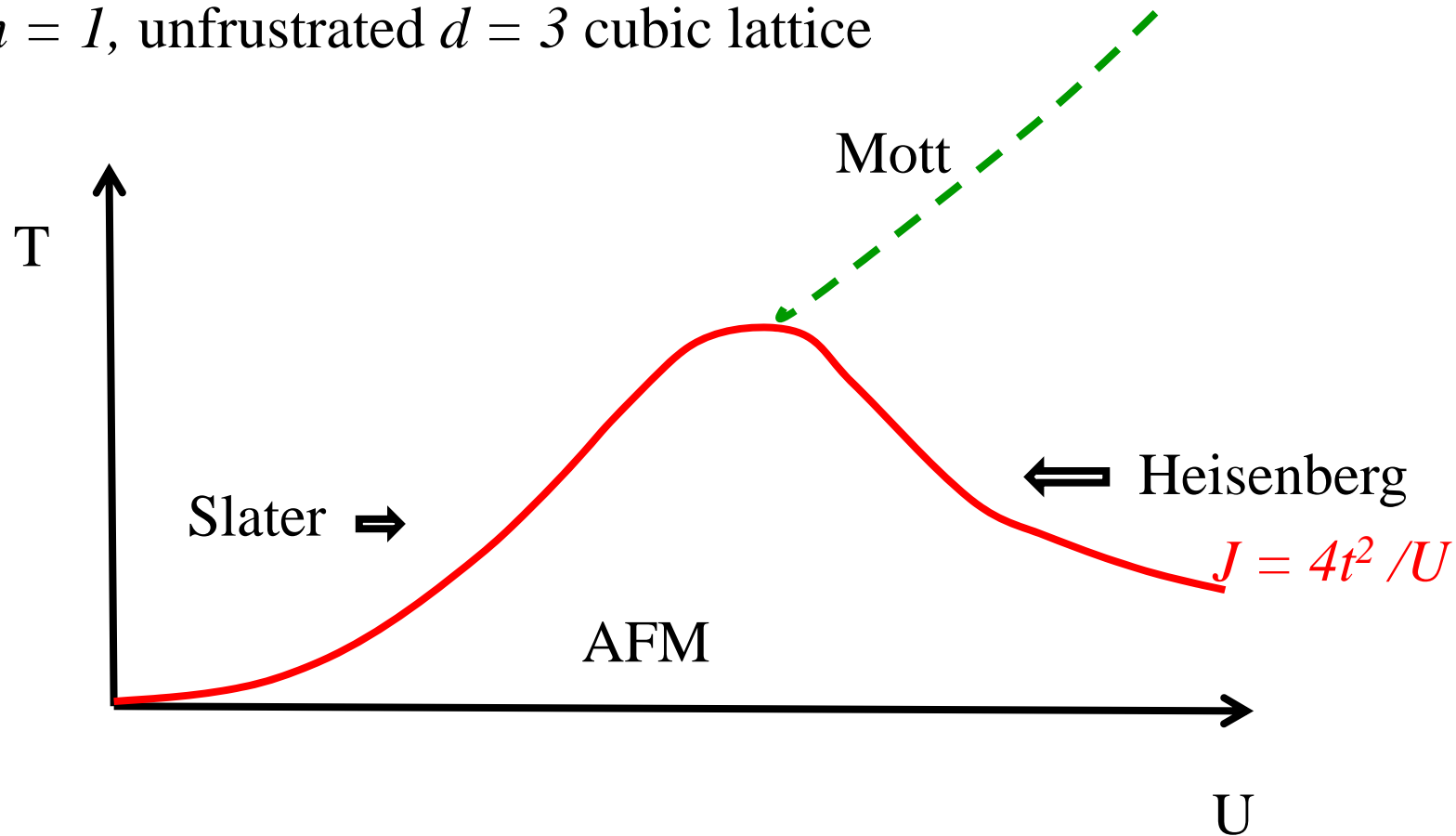


Effective model, Heisenberg: $J = 4t^2 / U$



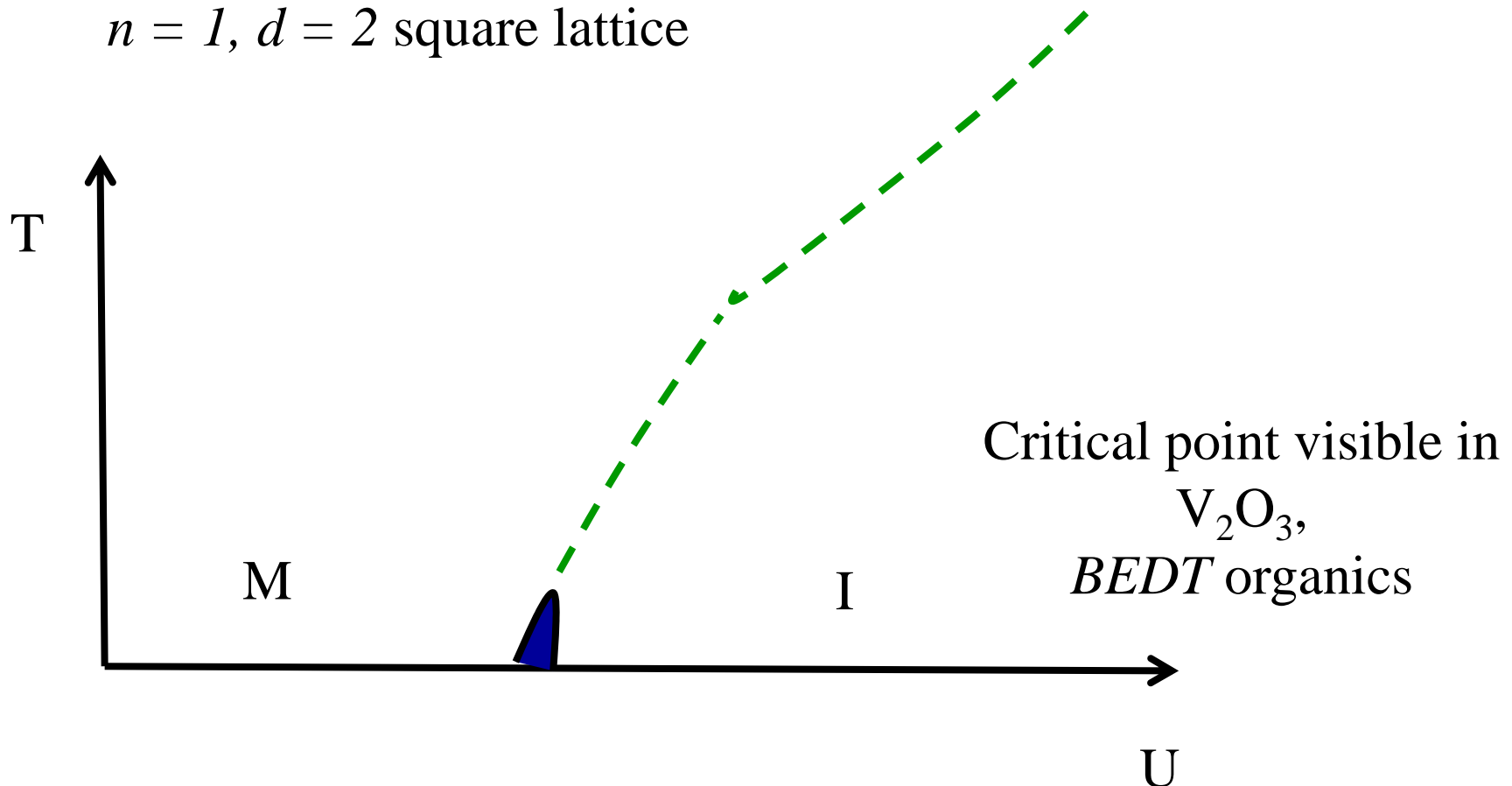
Weak-strong coupling, and Mott transition

$n = 1$, unfrustrated $d = 3$ cubic lattice



Local moment and Mott transition

$n = 1, d = 2$ square lattice



Understanding finite temperature phase from a *mean-field theory* down to $T = 0$



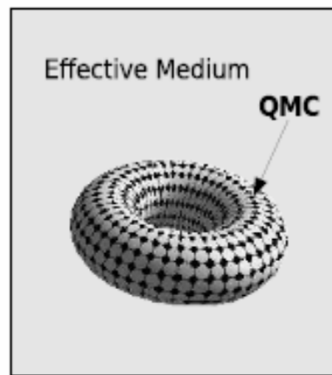
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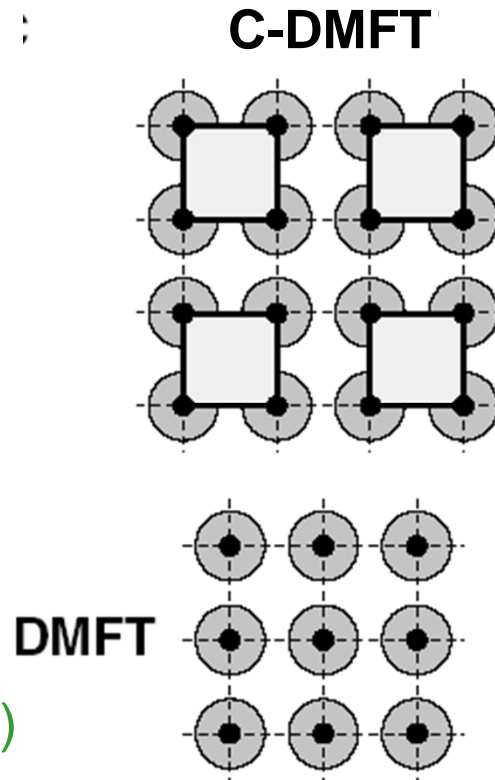


The method

2d Hubbard: Quantum cluster method



DCA



Hettler ...Jarrell...Krishnamurty PRB **58** (1998)

Kotliar et al. PRL **87** (2001)

M. Potthoff *et al.* PRL **91**, 206402 (2003).

REVIEWS

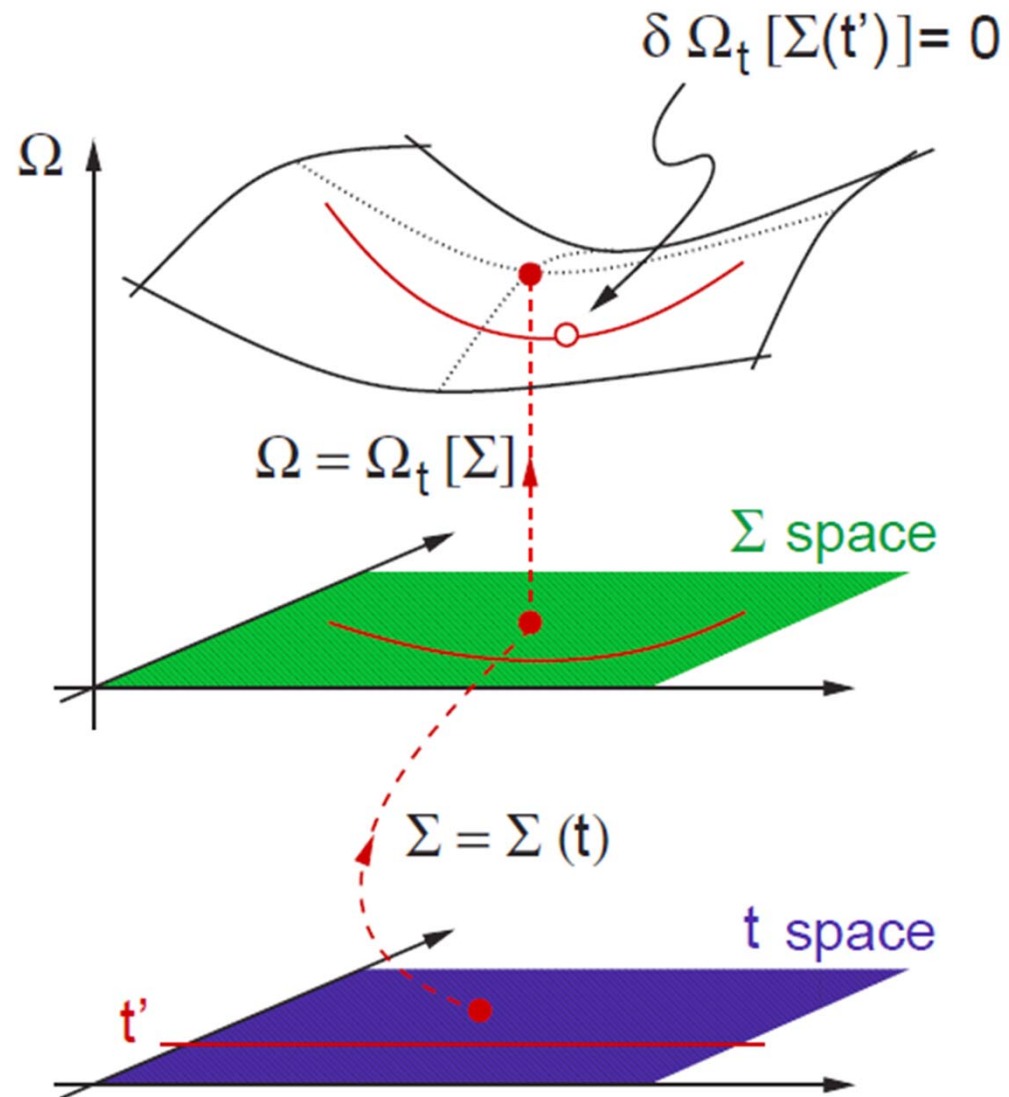
Maier, Jarrell et al., RMP. (2005)

Kotliar *et al.* RMP (2006)

AMST *et al.* LTP (2006)



DMFT as a stationary point



Another way to look at this (Potthoff)

$$\Omega_t[G] = \Phi[G] - \text{Tr}[(G_{0t}^{-1} - G^{-1})G] + \text{Tr} \ln(-G)$$

$$\frac{\delta\Phi[G]}{\delta G} = \Sigma$$

$$\Omega_t[\Sigma] = \Phi[G] - \text{Tr}[\Sigma G] - \text{Tr} \ln(-G_{0t}^{-1} + \Sigma)$$

Still stationary (chain rule)

$$\Omega_t[\Sigma] = F[\Sigma] - \text{Tr} \ln(-G_{0t}^{-1} + \Sigma)$$

SFT : Self-energy Functional Theory

With $F[\Sigma]$ Legendre transform of Luttinger-Ward funct.

$$\Omega_t[\Sigma] = F[\Sigma] + \text{Tr} \ln(-(G_0^{-1} - \Sigma)^{-1})$$

is stationary with respect to Σ and equal to grand potential there.

$$\Omega_t[\Sigma] = \Omega_{t'}[\Sigma] - \text{Tr} \ln(-(G_0'^{-1} - \Sigma)^{-1}) + \text{Tr} \ln(-(G_0^{-1} - \Sigma)^{-1}).$$

Vary with respect to parameters of the cluster (including Weiss fields)

Variation of the self-energy, through parameters in $H_0(\mathbf{t}')$

+ and -

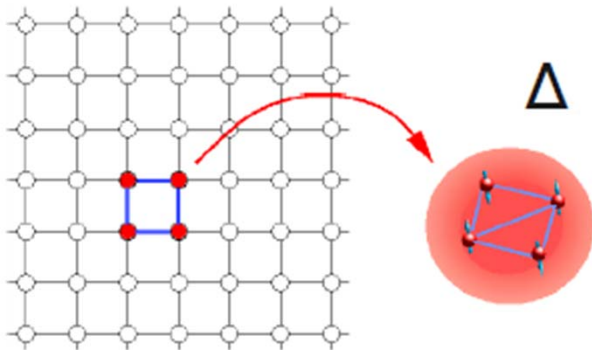
- Long range order:
 - Allow symmetry breaking in the bath (mean-field)
- Included:
 - Short-range dynamical and spatial correlations
- Missing:
 - Long wavelength p-h and p-p fluctuations



Solvers for the cluster-in-a-bath problem



C-DMFT

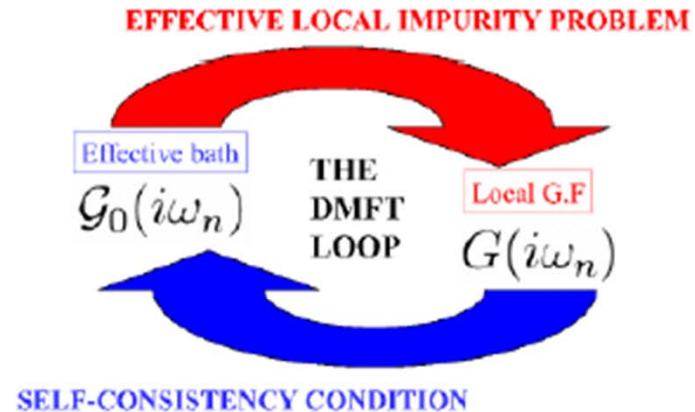


Mean-field is not a trivial problem! Many impurity solvers.

Here: continuous time QMC

P. Werner, PRL 2006
 P. Werner, PRB 2007
 K. Haule, PRB 2007

$$Z = \int \mathcal{D}[\psi^\dagger, \psi] e^{-S_c - \int_0^\beta d\tau \int_0^\beta d\tau' \sum_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger(\tau) \Delta_{\mathbf{k}}(\tau, \tau') \psi_{\mathbf{k}}(\tau')}$$



$$\Delta(i\omega_n) = i\omega_n + \mu - \Sigma_c(i\omega_n)$$

$$= \left[\sum_{\tilde{\mathbf{k}}} \frac{1}{i\omega_n + \mu - t_c(\tilde{\mathbf{k}}) - \Sigma_c(i\omega_n)} \right]^{-1}$$

At finite T , solving cluster in a bath problem

- Continuous-time Quantum Monte Carlo calculations to sum all diagrams generated from expansion in powers of hybridization.
 - P. Werner, A. Comanac, L. de' Medici, M. Troyer, and A. J. Millis, Phys. Rev. Lett. **97**, 076405 (2006).
 - K. Haule, Phys. Rev. B **75**, 155113 (2007).



Outline

- Pseudogap
 - From AFM fluctuations
- Weak vs strong coupling
- Dynamical mean-field theory: crossover
- **Phase diagram and pseudogap at strong coupling**





Giovanni Sordi

Finite T phase diagram

CTQMC as solver for cluster in a bath ($T=0$).



Patrick Sémon

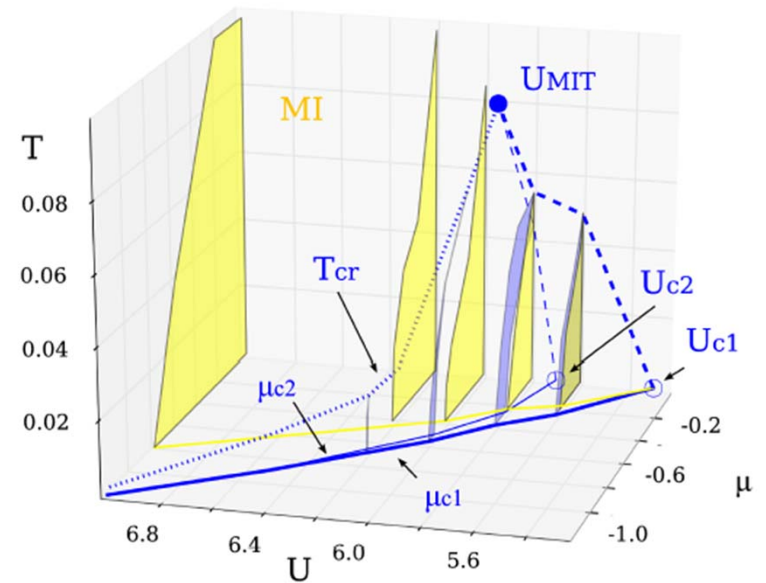
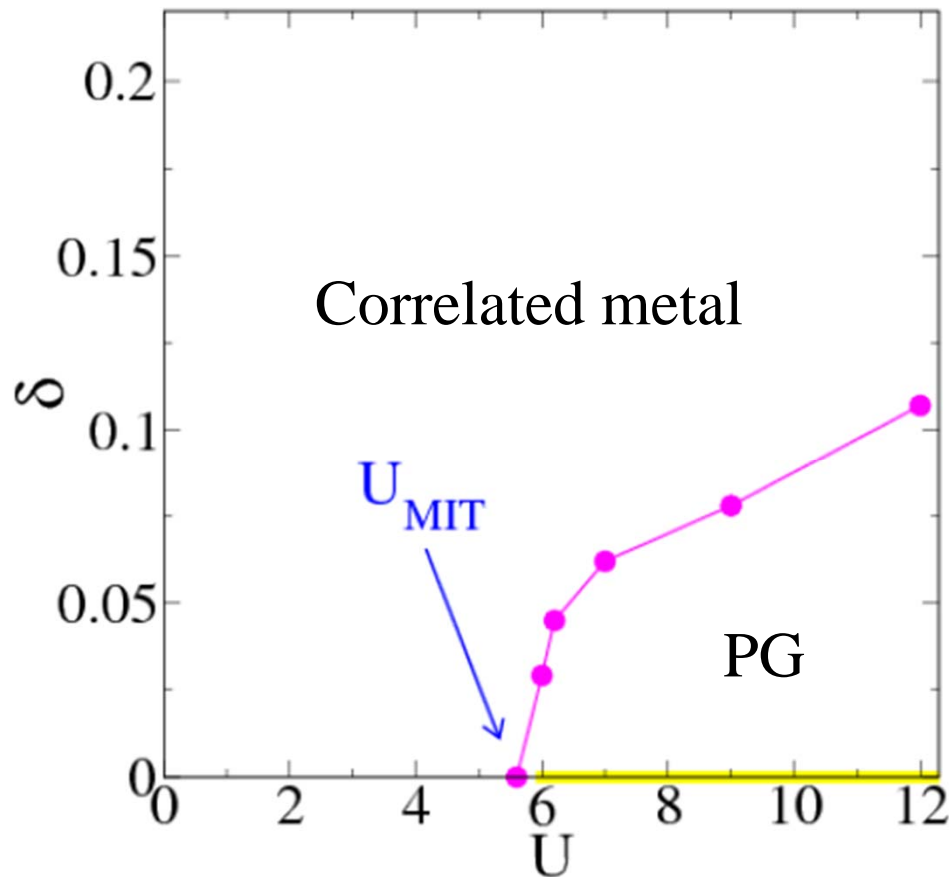


Kristjan Haule

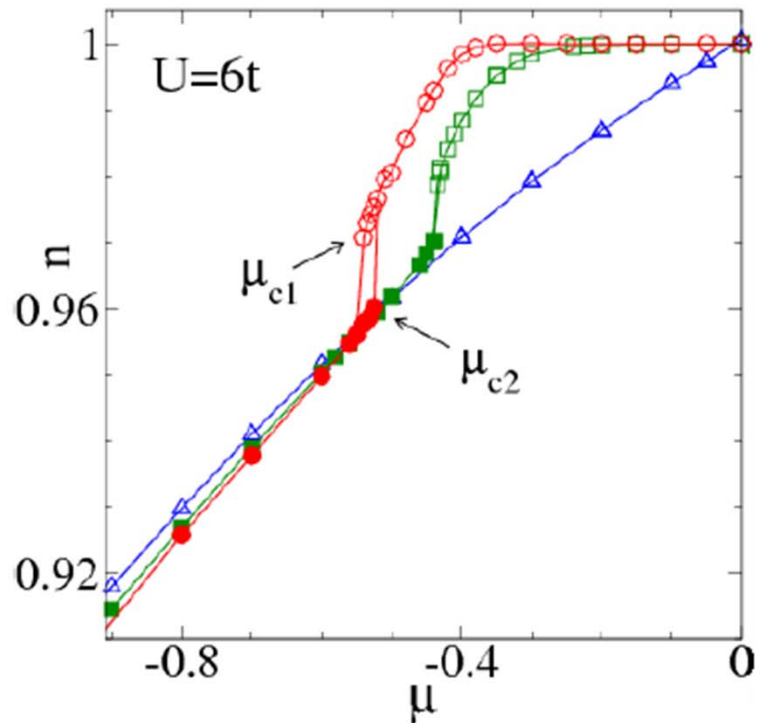


Link to Mott transition up to optimal doping

Doping dependence of critical point as a function of U



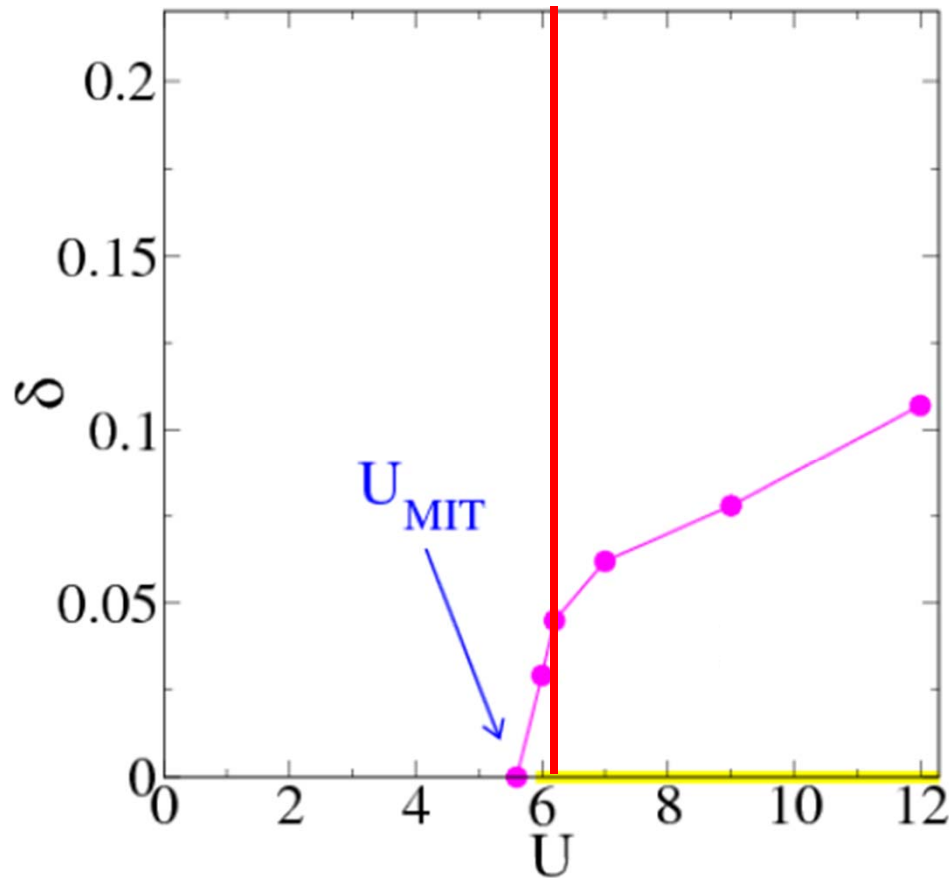
First order transition at finite doping



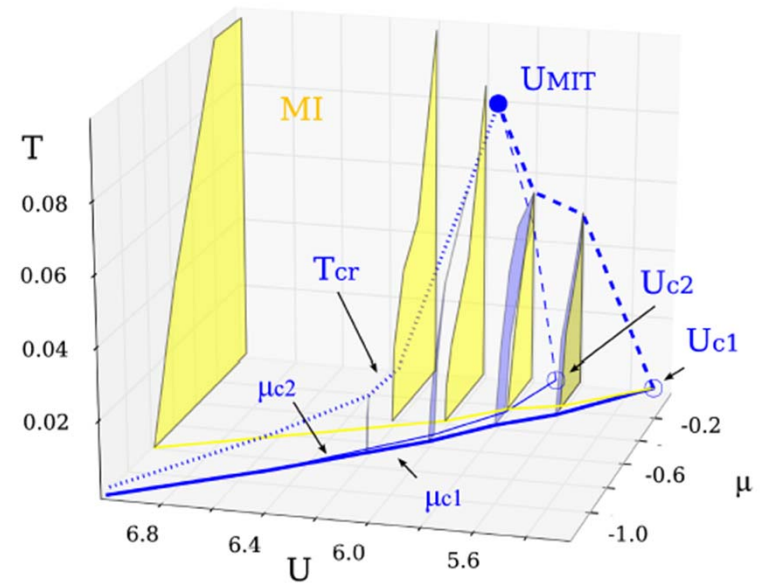
$n(\mu)$ for several temperatures:
 $T/t = 1/10, 1/25, 1/50$

Link to Mott transition up to optimal doping

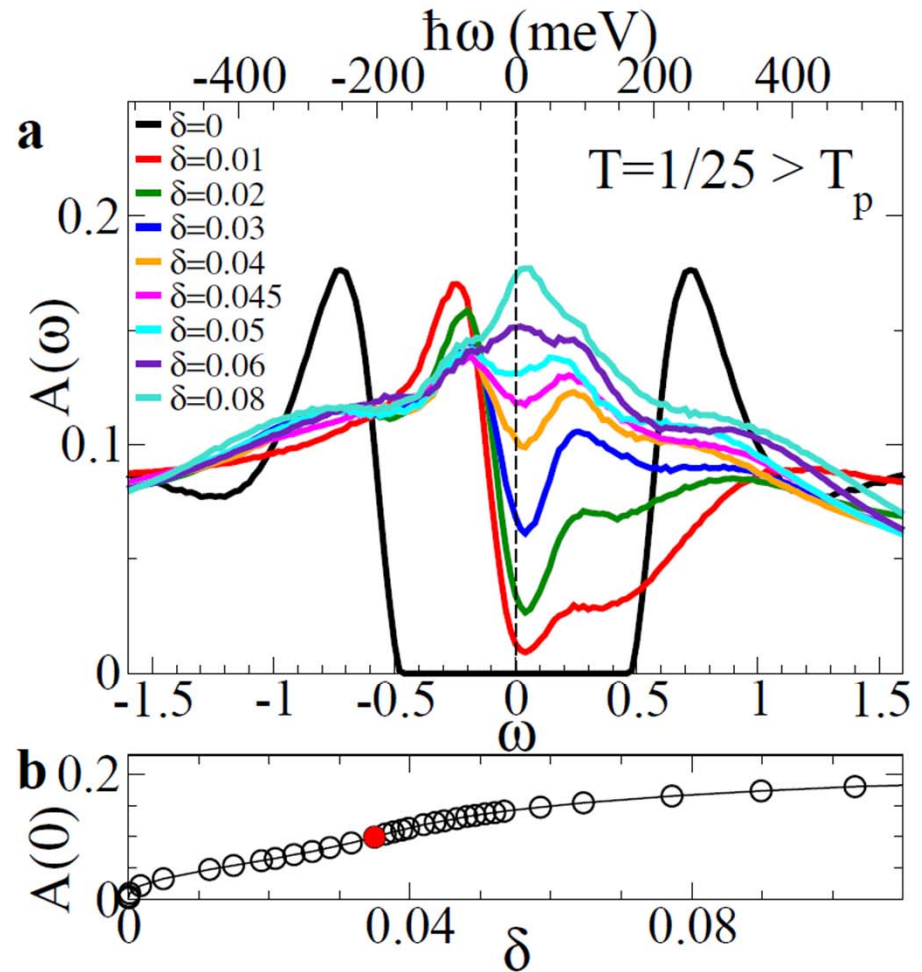
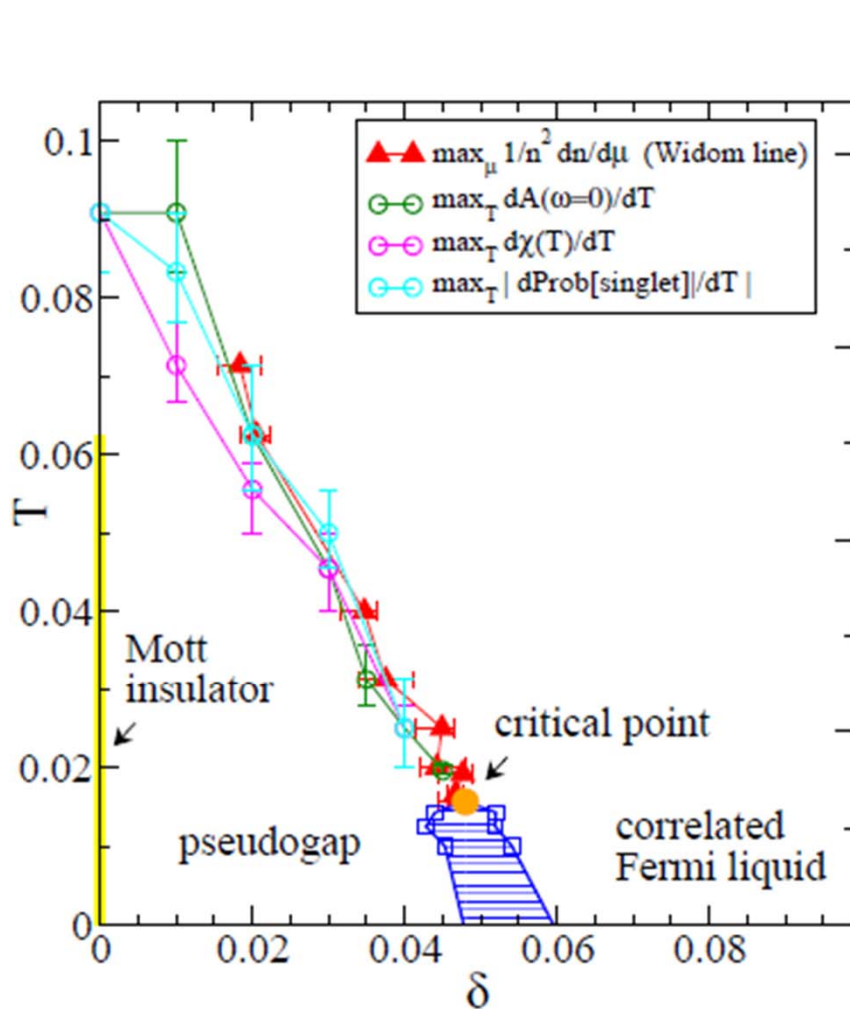
Doping dependence of critical point as a function of U



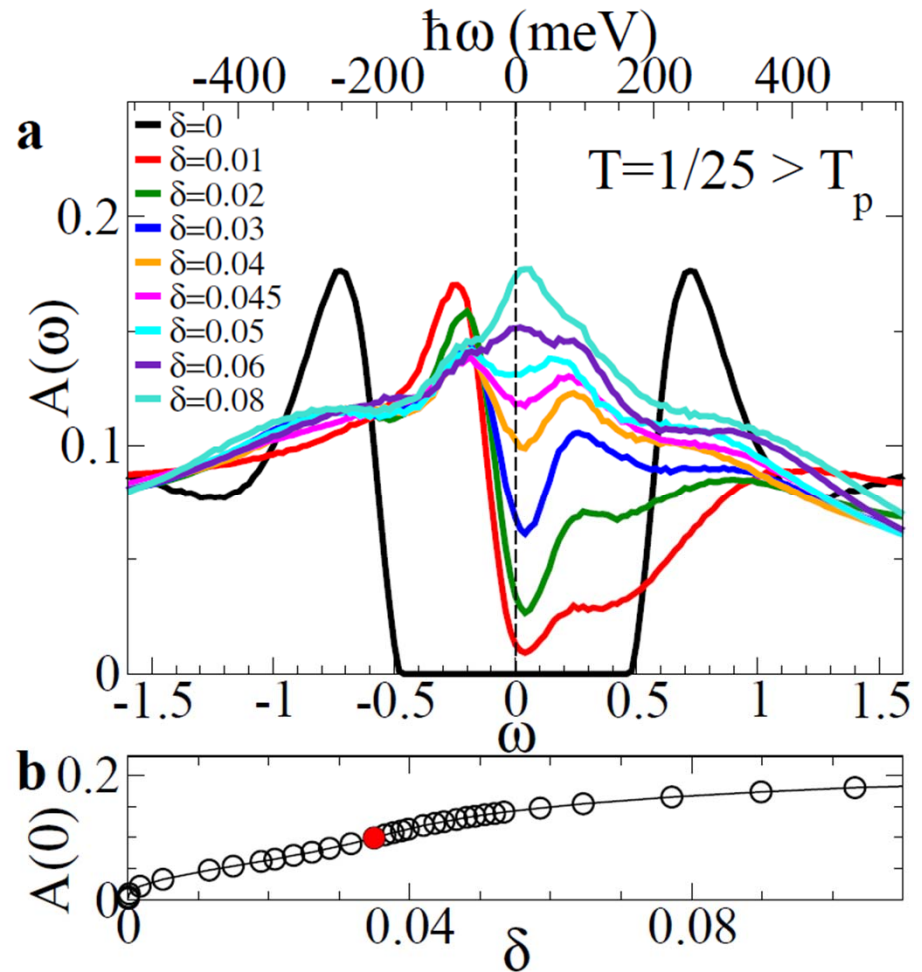
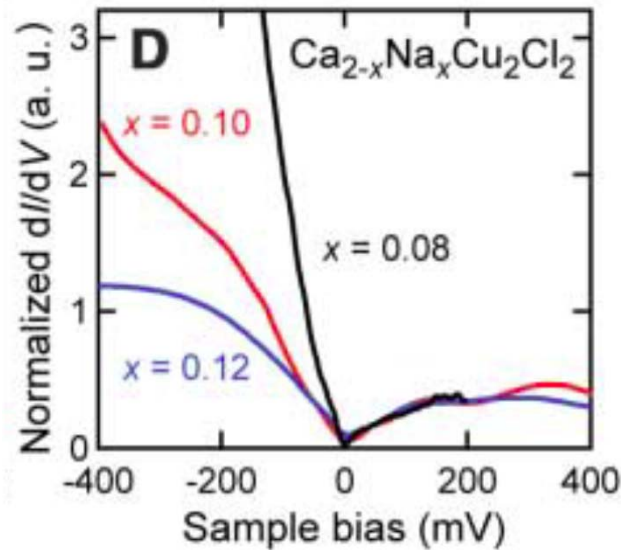
Smaller D and S



Density of states



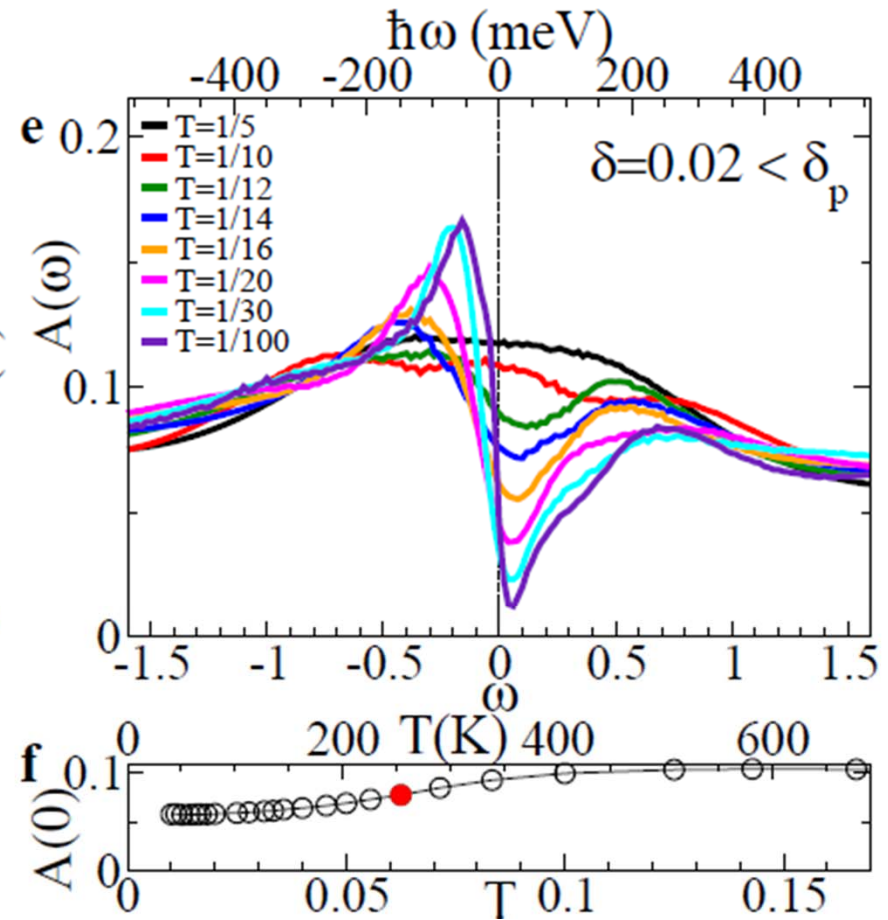
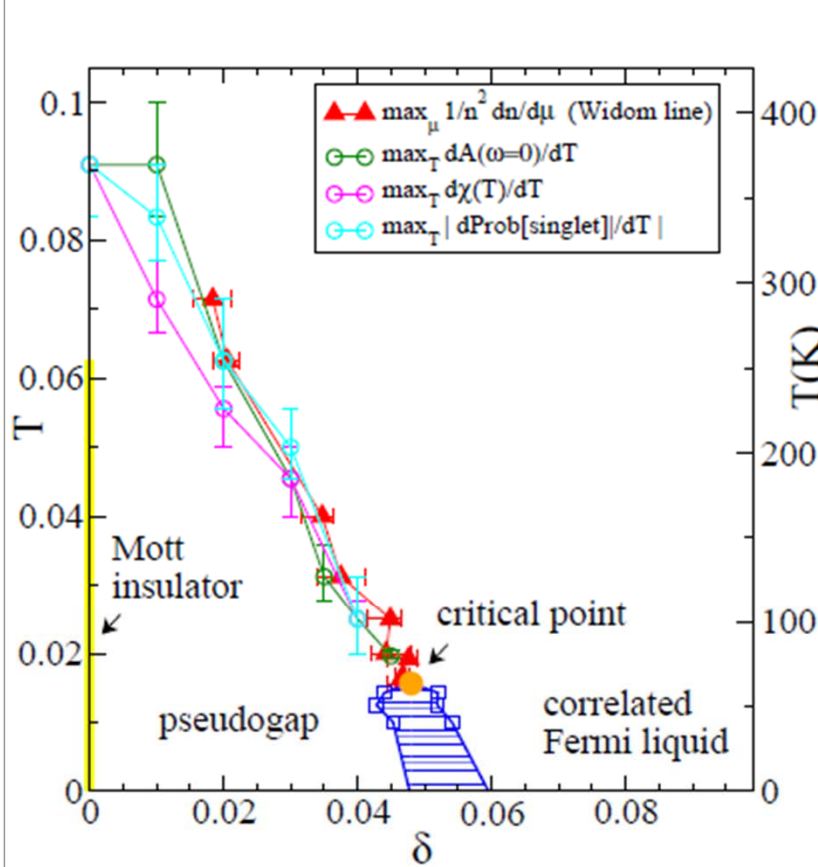
Density of states



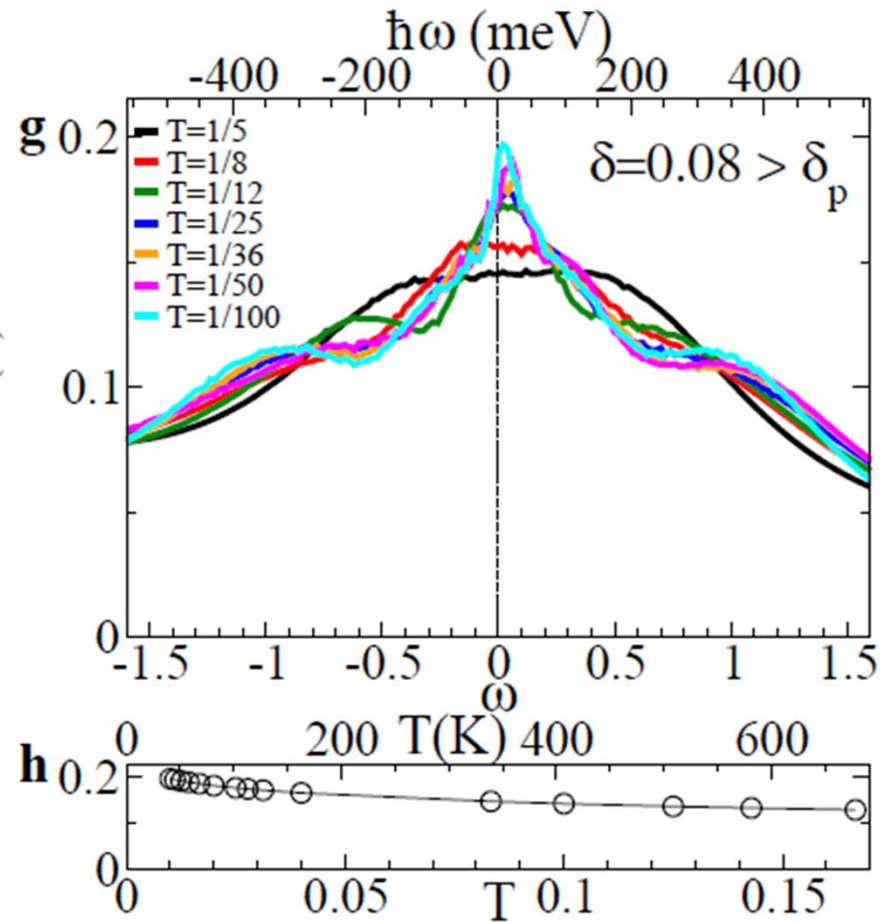
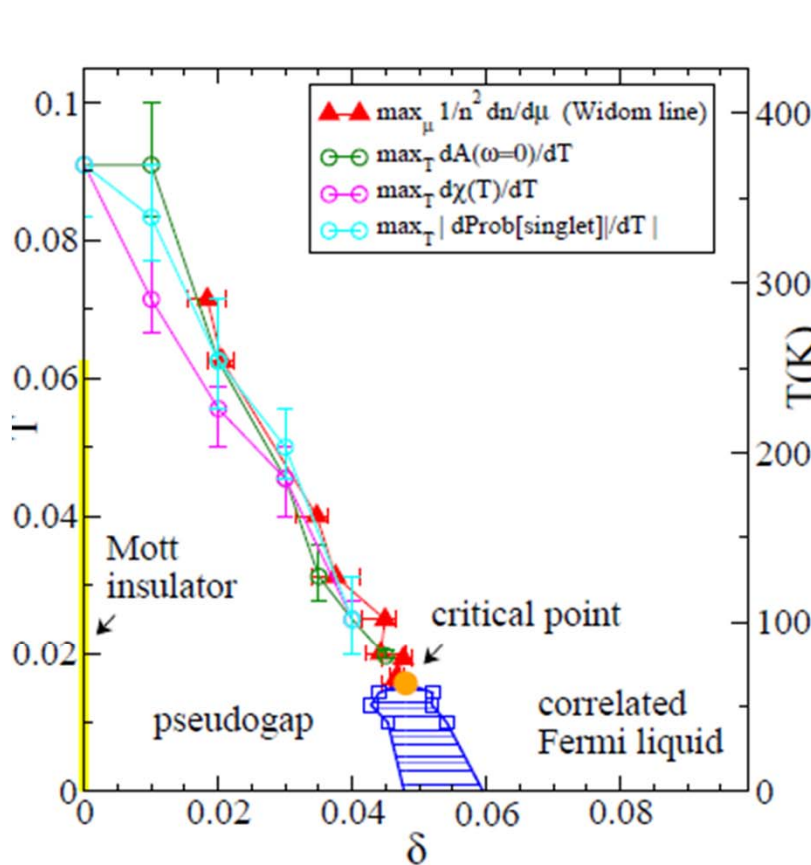
Khosaka et al. *Science* **315**, 1380 (2007);



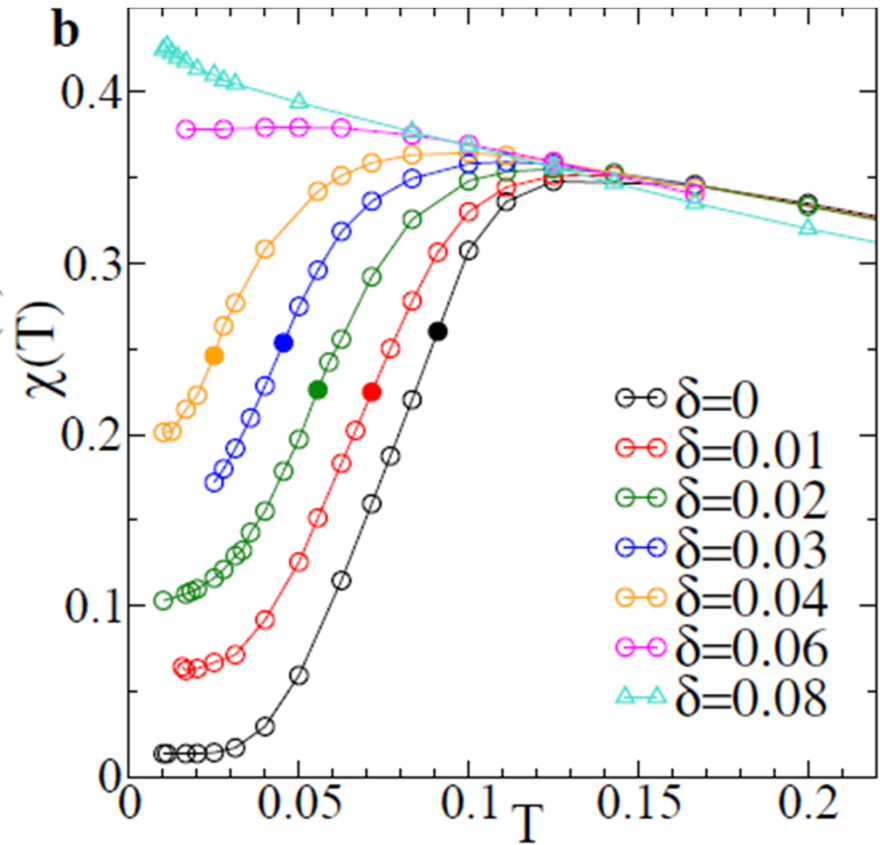
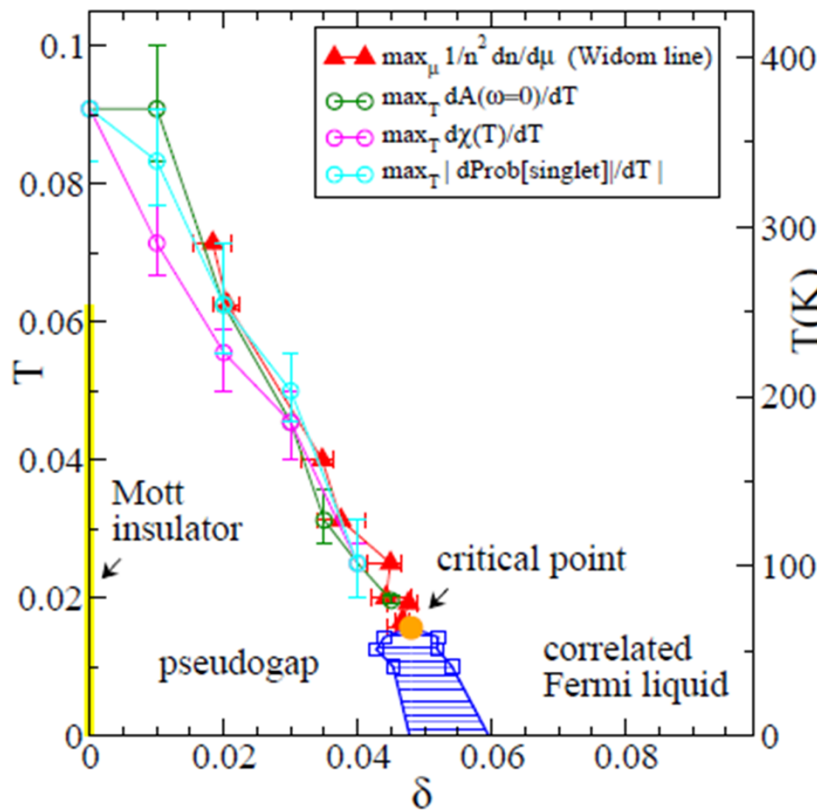
Density of states



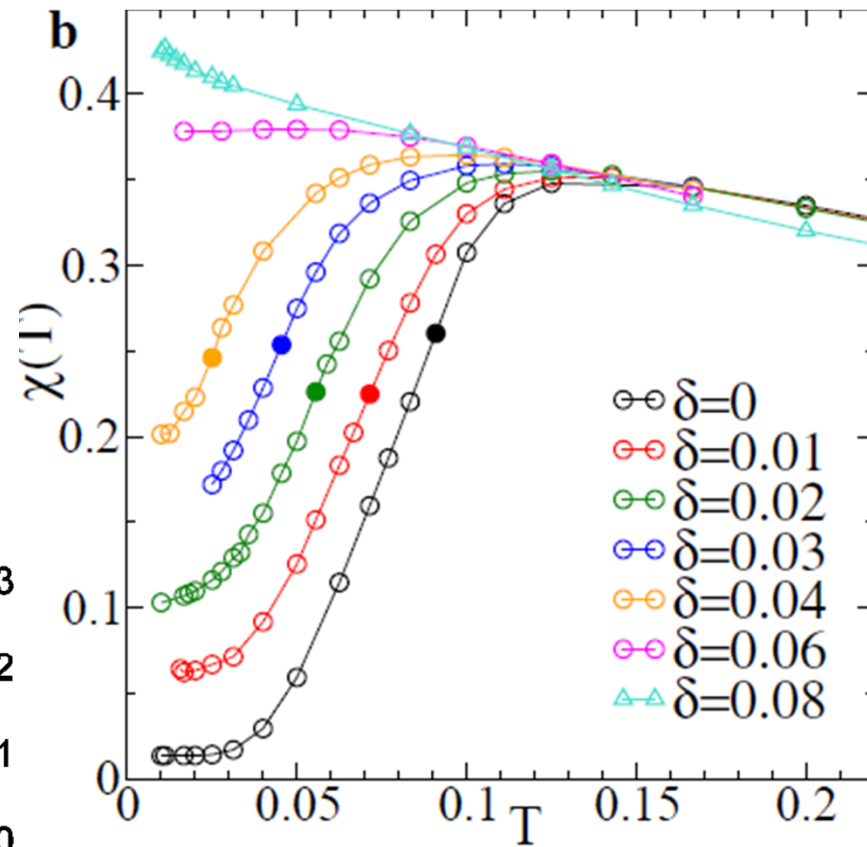
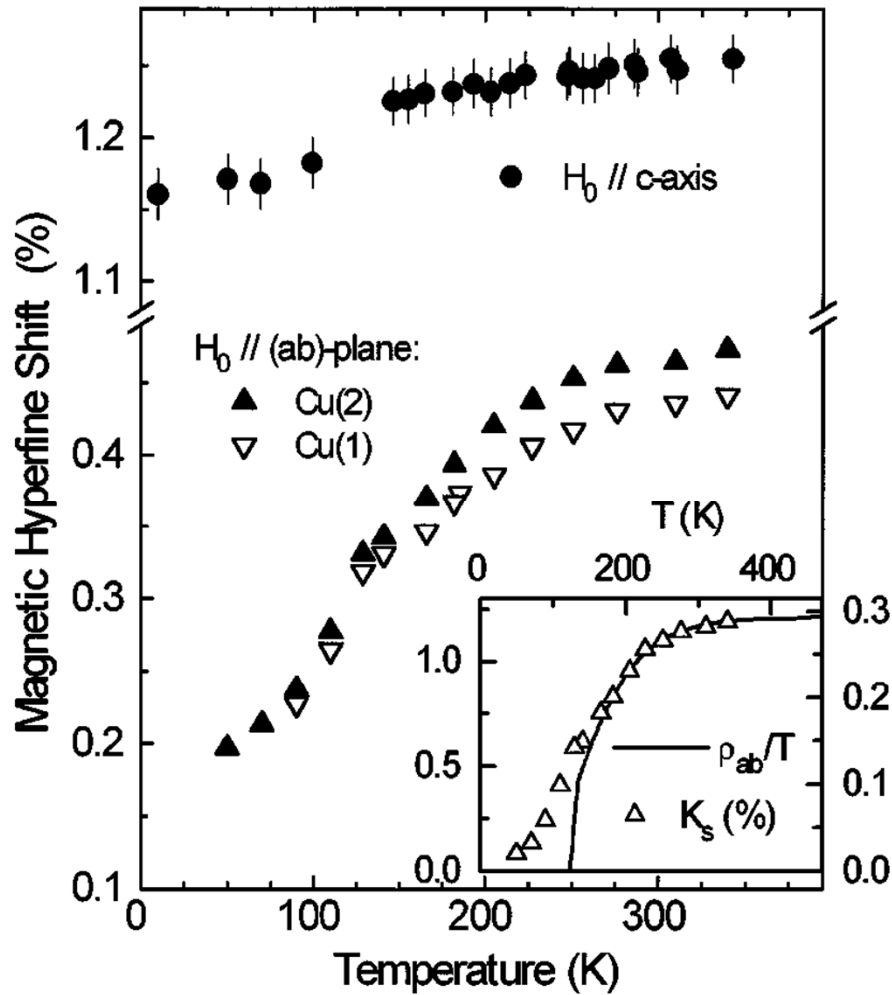
Density of states



Spin susceptibility



Spin susceptibility



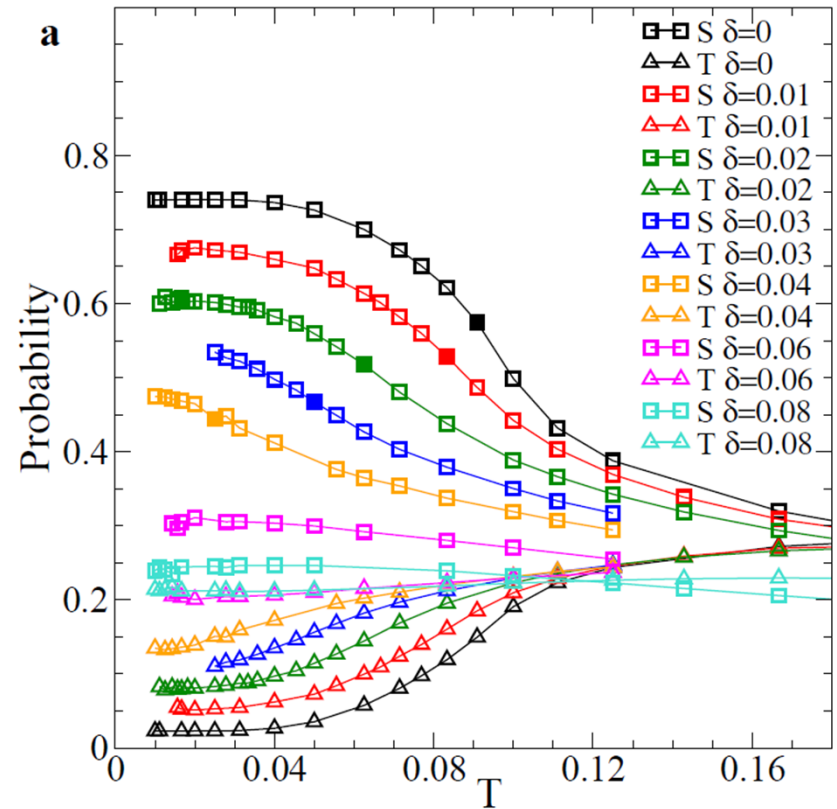
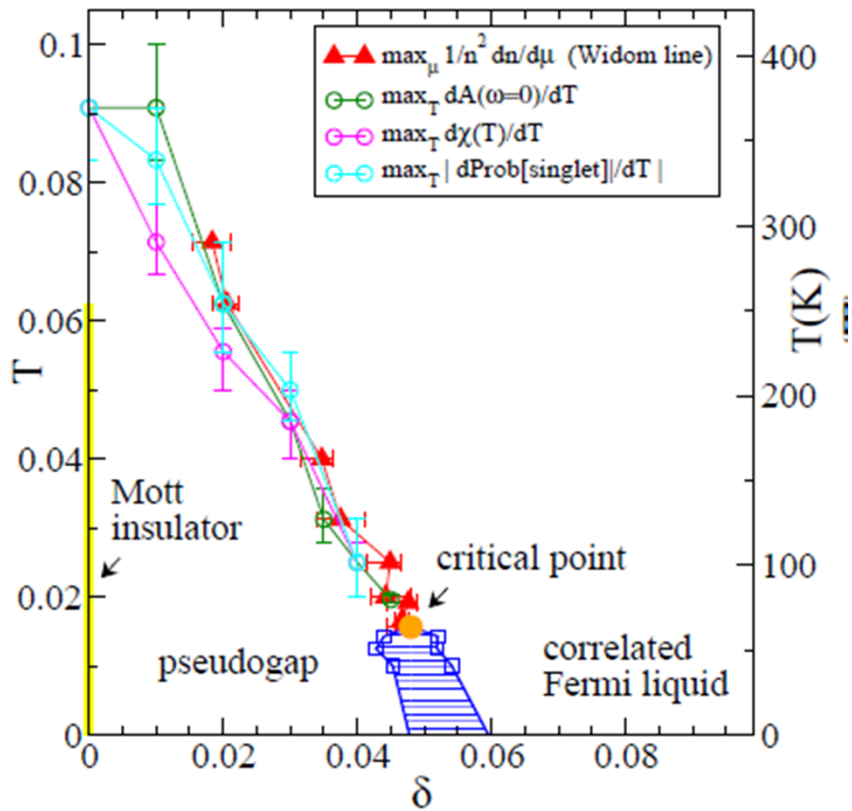
Underdoped Hg1223

Julien et al. PRL **76**, 4238 (1996)

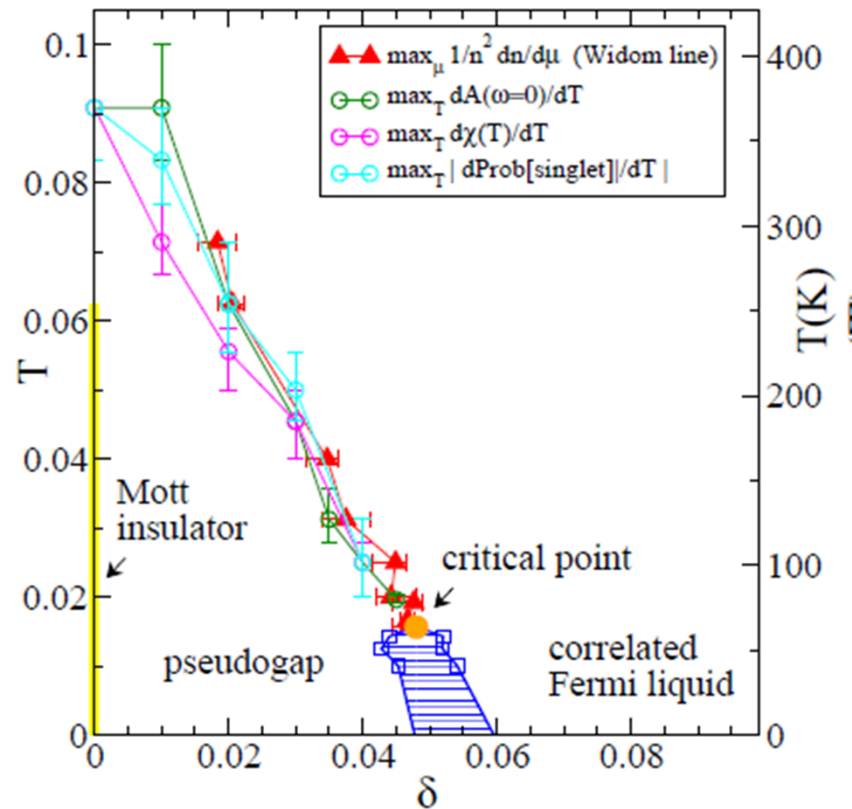


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Plaquette eigenstates



Pseudogap T^* along the Widom line





Giovanni Sordi



Patrick Sémon



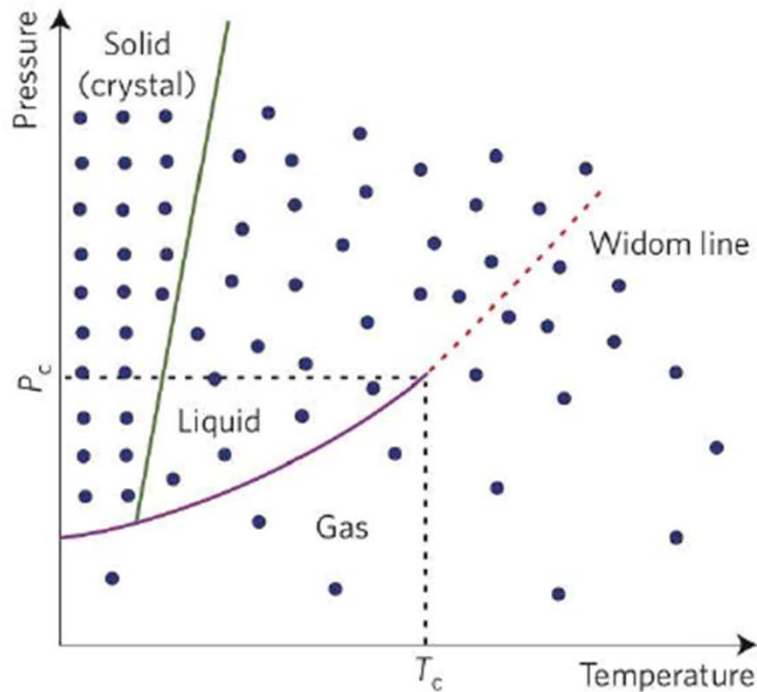
Kristjan Haule

The Wisdom line

G. Sordi, *et al.* Scientific Reports 2, 547 (2012)



What is the Widom line?

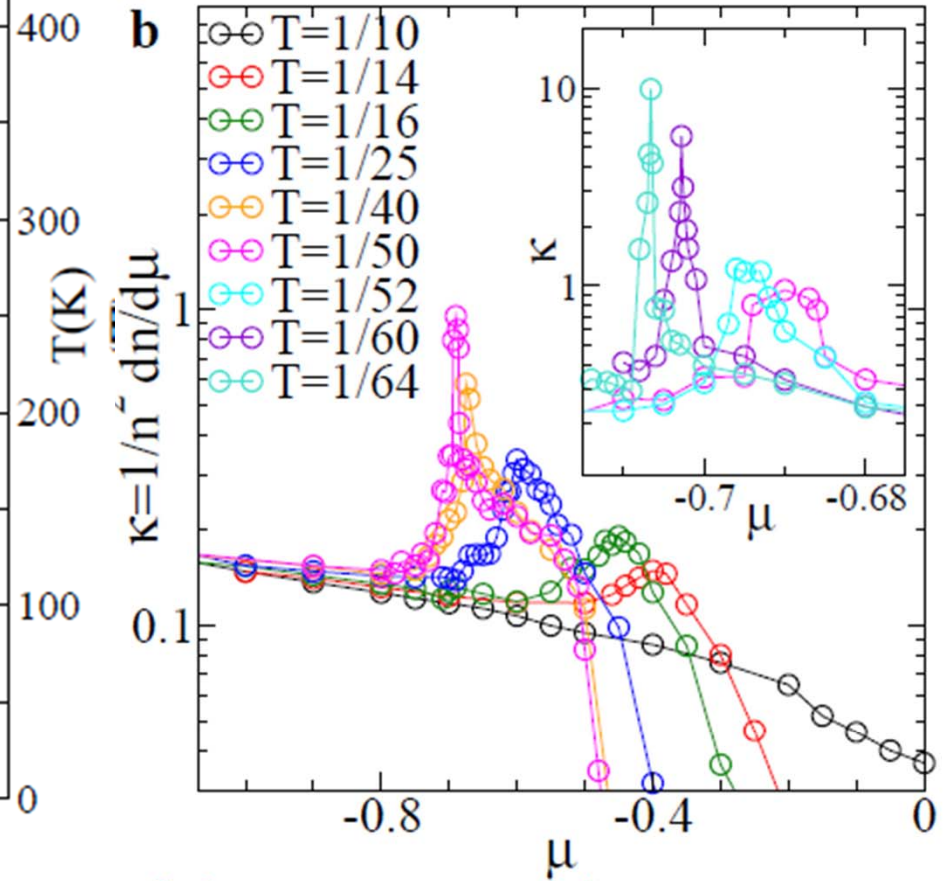
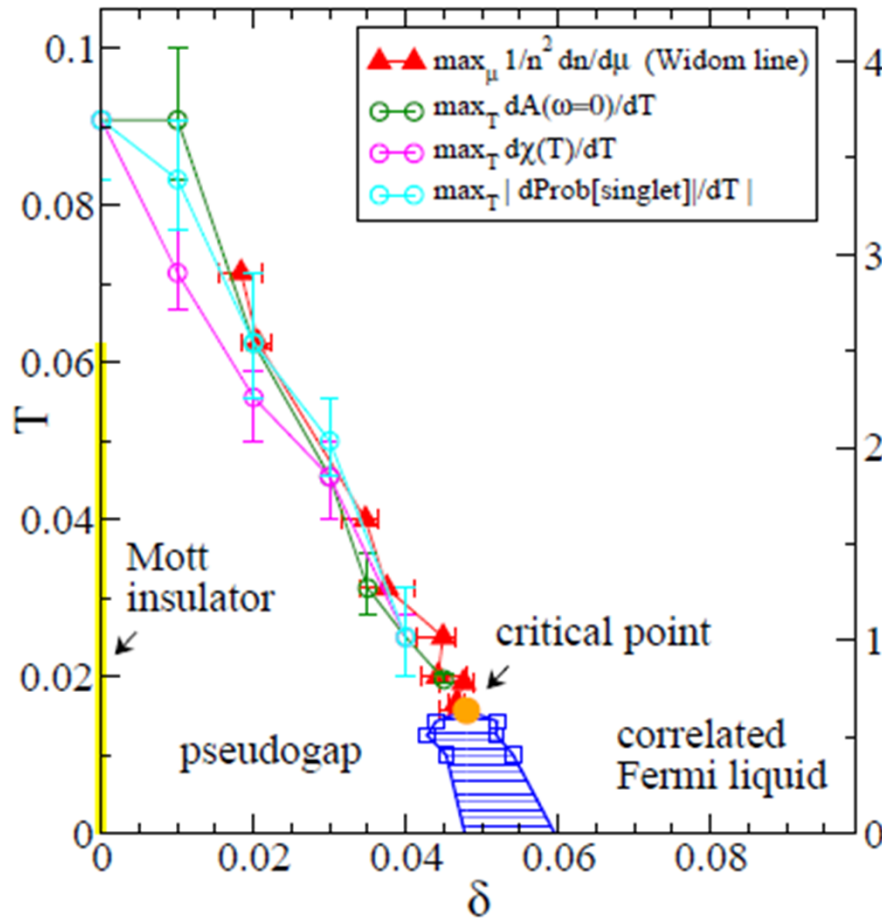


McMillan and Stanley, Nat Phys 2010

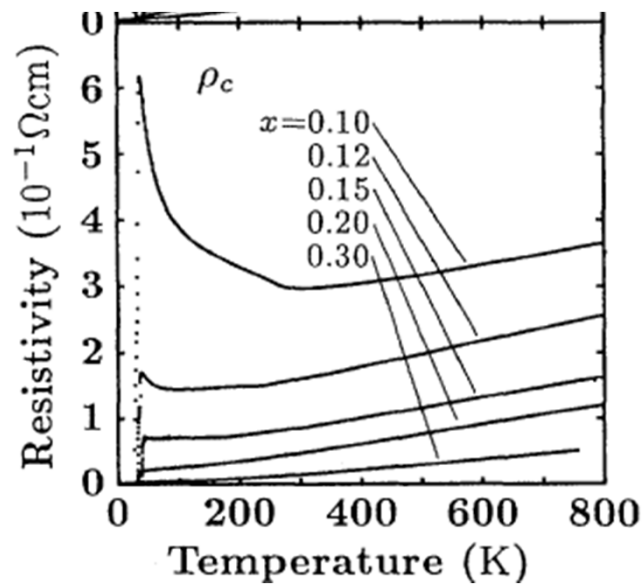
- ▶ it is the continuation of the coexistence line in the supercritical region
- ▶ line where the **maxima of different response functions** touch each other asymptotically as $T \rightarrow T_p$
- ▶ liquid-gas transition in water: max in isobaric heat capacity C_p , isothermal compressibility, isobaric heat expansion, etc
- ▶ **DYNAMIC crossover arises from crossing the Widom line!**
water: Xu et al, PNAS 2005, Simeoni et al Nat Phys 2010



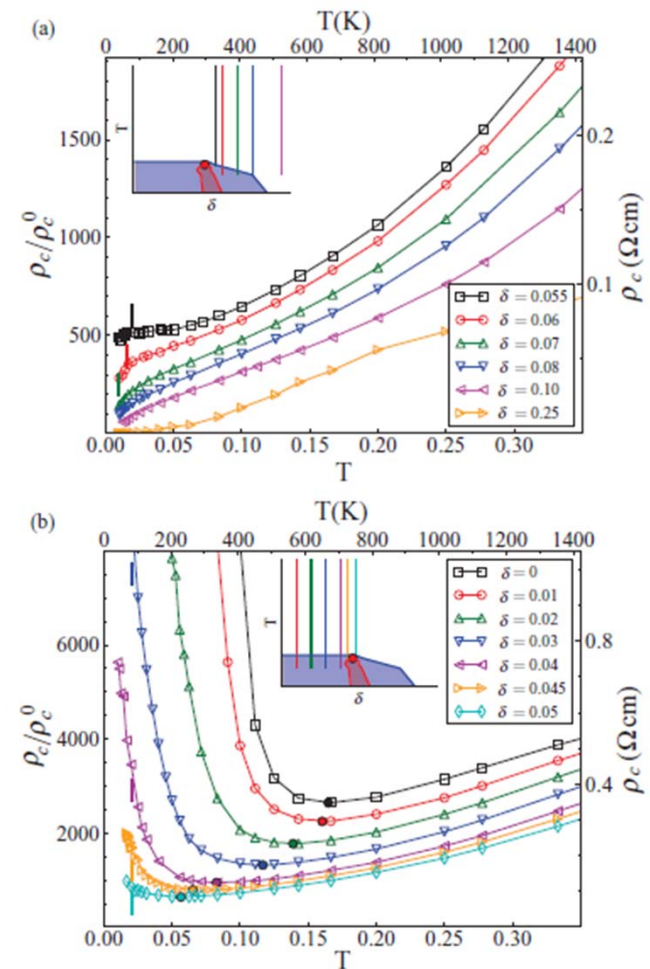
Phase diagram



C-axis resistivity



Y. Nakamura, S. Uchida, PRB **47**, 8369 (1993)



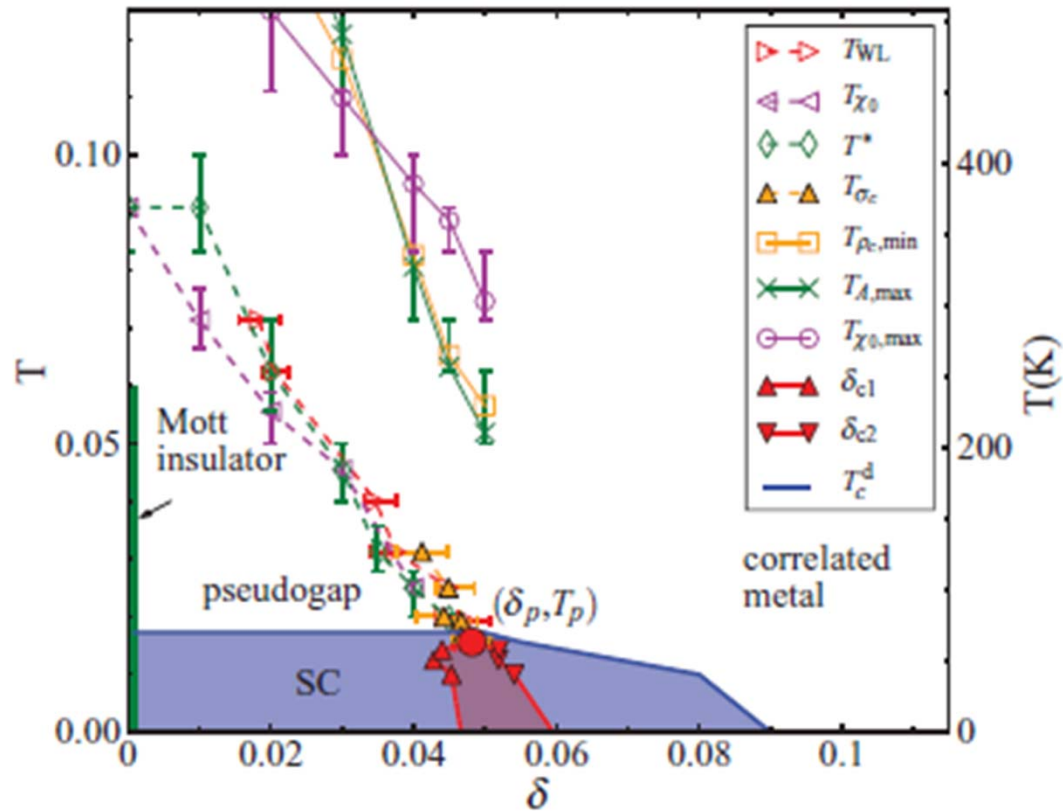
G. Sordi, P. Sémon, K. Haule, A.-M.S.T

PRB, **87**, 041101(R) (2013)



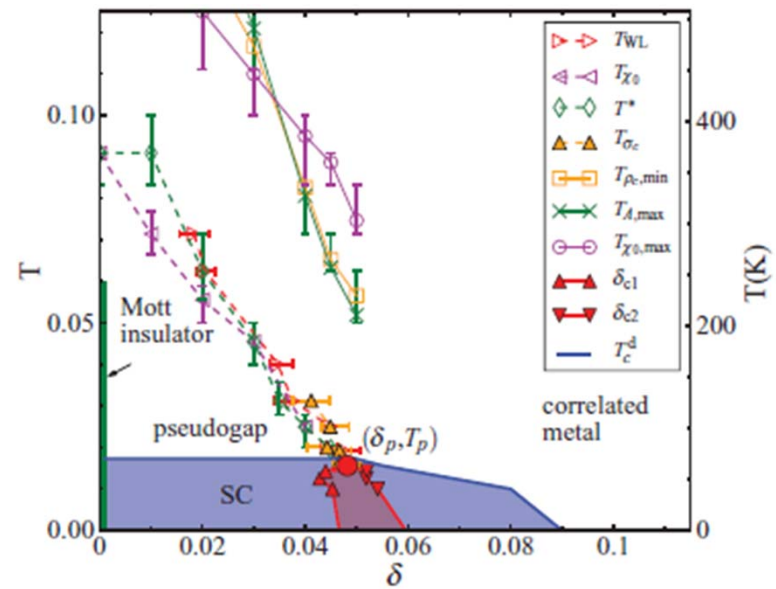
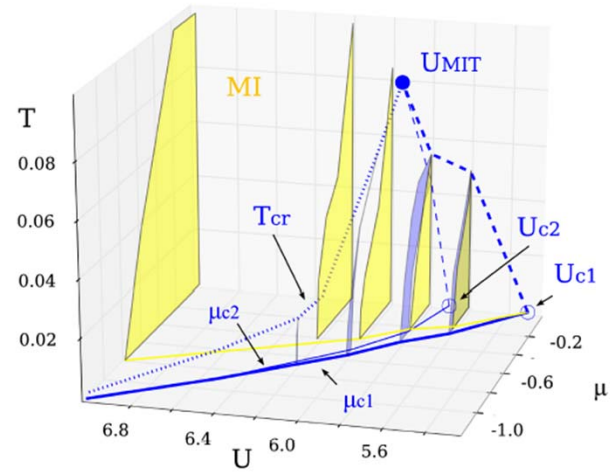
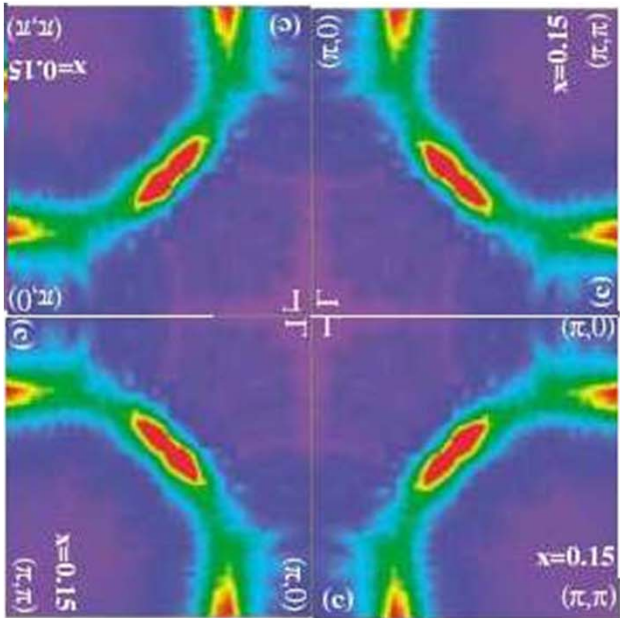
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Two crossover lines



G. Sordi, P. Sémon, K. Haule, A.-M.S.T
 PRB, **87**, 041101(R) (2013)

Summary



Main collaborators



Giovanni Sordi



Kristjan Haule



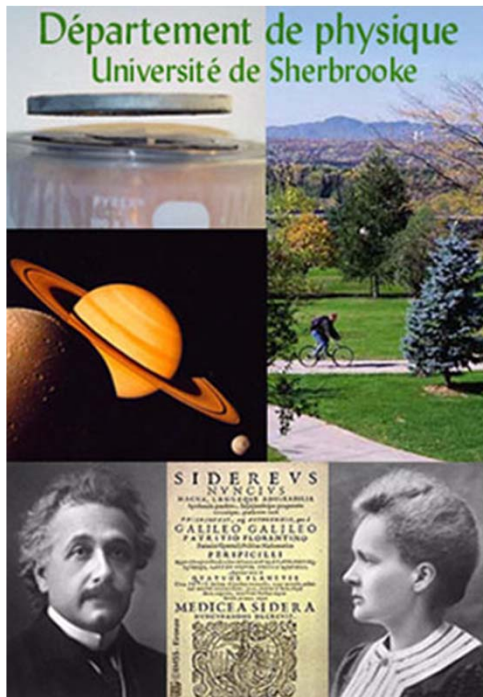
Patrick Sémon



David Sénéchal



André-Marie Tremblay



Le regroupement québécois sur les matériaux de pointe



Sponsors:



Mammoth



 **compute + calcul**
CANADA

High Performance Computing

CREATING KNOWLEDGE
DRIVING INNOVATION
BUILDING THE DIGITAL ECONOMY

Le calcul de haute performance

CRÉER LE SAVOIR
ALIMENTER L'INNOVATION
BÂTIR L'ÉCONOMIE NUMÉRIQUE


Calcul Québec


Canada Foundation for Innovation
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Merci

Thank you